

*A depth-averaged, two-phase flow code for hazard mapping that satisfies both hydraulic and granular flow extremes.*

By

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# The Titan project

## Aims

- Create realistic 2D & 3D models that simulate geophysical mass flow
  - Integrate and communicate information from several linked sources including:
    - .- Simulation results
    - .- Remote sensing data
    - .- GIS data
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# Current applications

- One phase dense granular flows

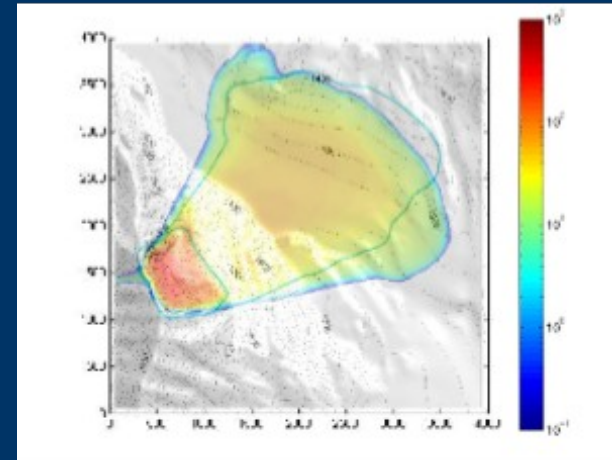
Rock avalanches  
Block-and-ash flows

- Two phase dense flows

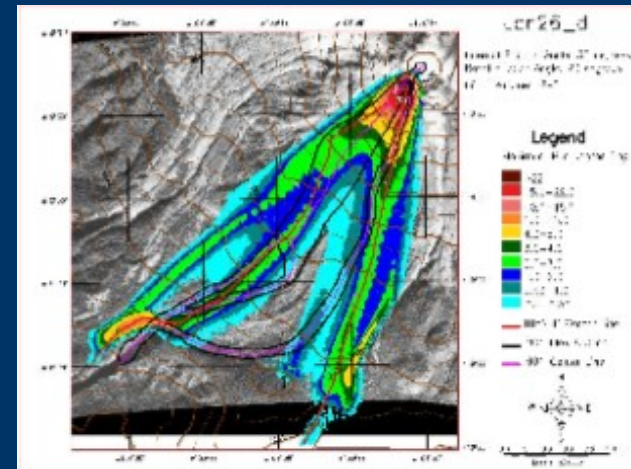
Debris flows

- Hazard maps & risk assessment

- Probabilistic analysis



Frank slide



Colima 1991

# Mathematical basis

## Savage-Hutter (1989) Model

- Provides a framework for later, more sophisticated models.

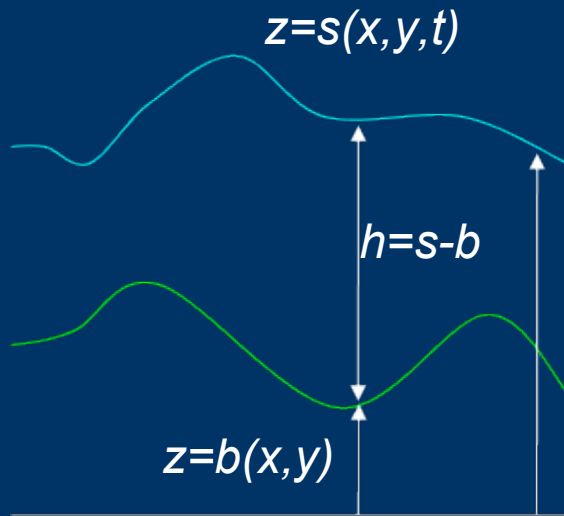
- The model is based on assumptions:

Aspect ratio of the flow is small

Granular flows have Coulomb friction behaviour

Top surface is stress-free

Top surface and bed surface represented as functions  $z=s(x,y,t)$  and  $z=b(x,y)$



The equations are depth averaged:

$$h \overline{v}_x = \int_b^s v_x dz, \quad h \overline{v}_y = \int_b^s v_y dz,$$

geophysical

mass

# Problems with the one phase model

- Several natural flows are multiphase
- Savage-Hutter model is only for dry avalanches
- Interstitial fluid modelled in Titan by reducing bed and internal friction angles

However,

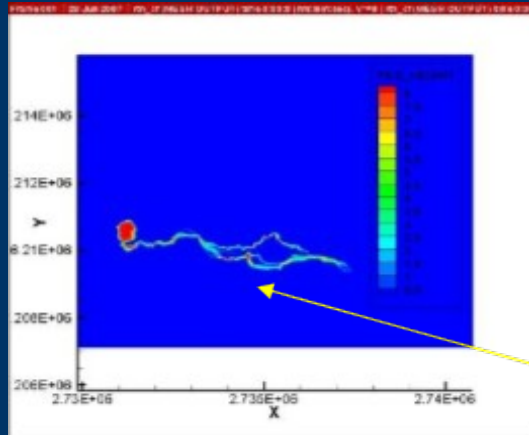
- Interstitial fluid can greatly alter the dynamics of the flow .
- Flows can achieve very high velocities, and can travel long distances.
- Water may be supplied by rainfall, melting of snow or ice or by overflow of a crater lake.



# Pitman and Le model

- *Pitman and Le, (2005)* expand the applications of TITAN to debris flows with various water contents
- Independent equations of motion specified for each phase, along with an equation for volume fractions
- Solid-Fluid interaction modelled using a drag force term

# Validation: actual prediction of the 2007 Mt Rapehu Lahar



Previously non recognized bifurcation predicted by Titan



Actual path



Bank levee

Titan prediction

## *Problems with Pitman model*

- Numerics break down at flow-front
- Non-physical velocities, as conserved variables, are divided by very small values
- It does not work well close to the extremes of pure dry or pure water flows.



# Cordoba, Pitman and Sheridan model

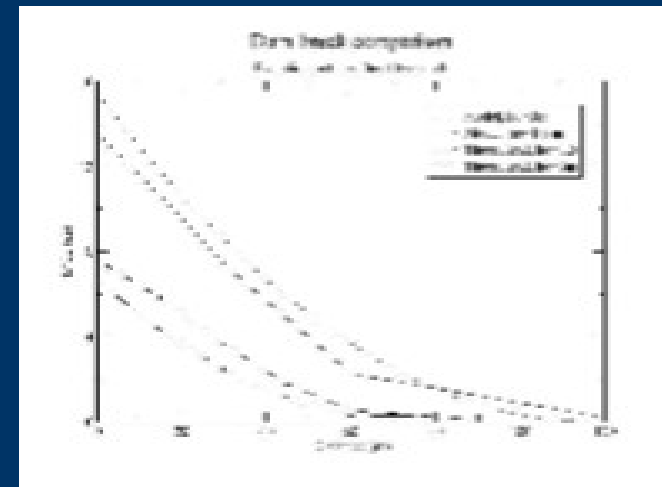
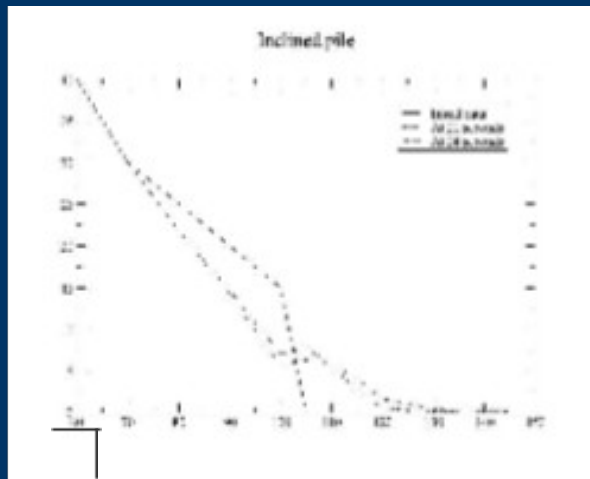
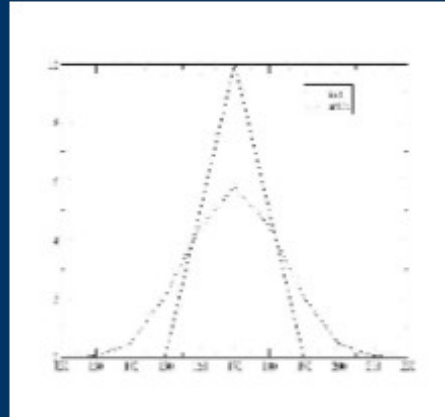
## Characteristics of the proposed model

- No limit on the fluid phase content
  - Pure dry avalanches modeled as a mixture of near maximum pack concentration of solids plus air.
  - Solid phase modeled as a granular flow that obeys the Coulomb friction law.
  - Fluid phase modeled as shallow water that follows the Darcy-Weisbach approach for friction.
  - Phase equations coupled through a drag coefficient and pressure.
  - Mathematically the equations become the Savage and Hutter (1889) approach for pure solids and to pure shallow water (e.g. see the model of Cun-hong et al., 2006) in case of zero solids concentration
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# Validation

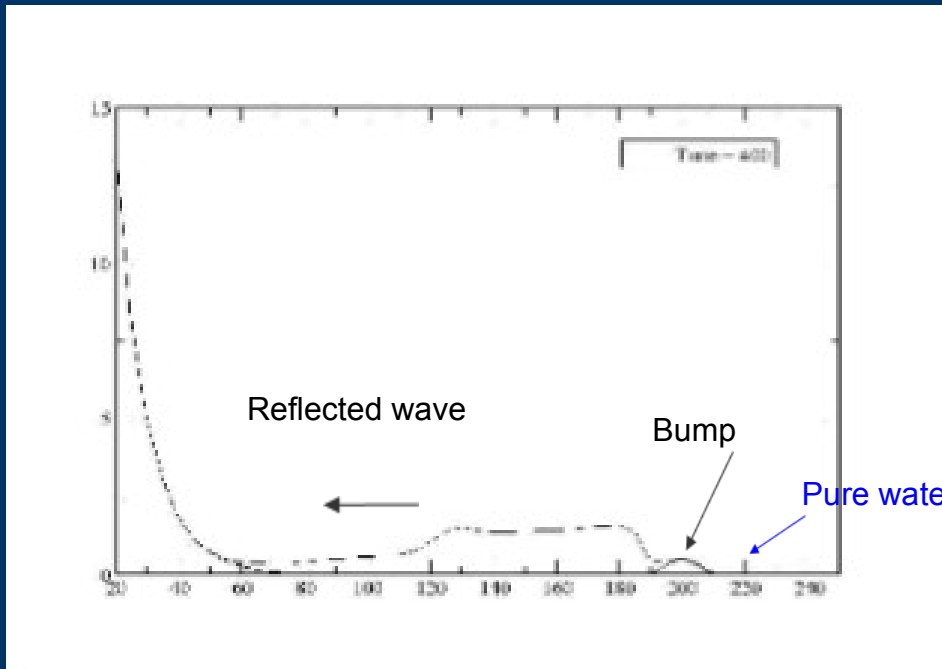
For purposes of validation and analytical testing, the one dimensional version of the equations was tested:

## Symmetry

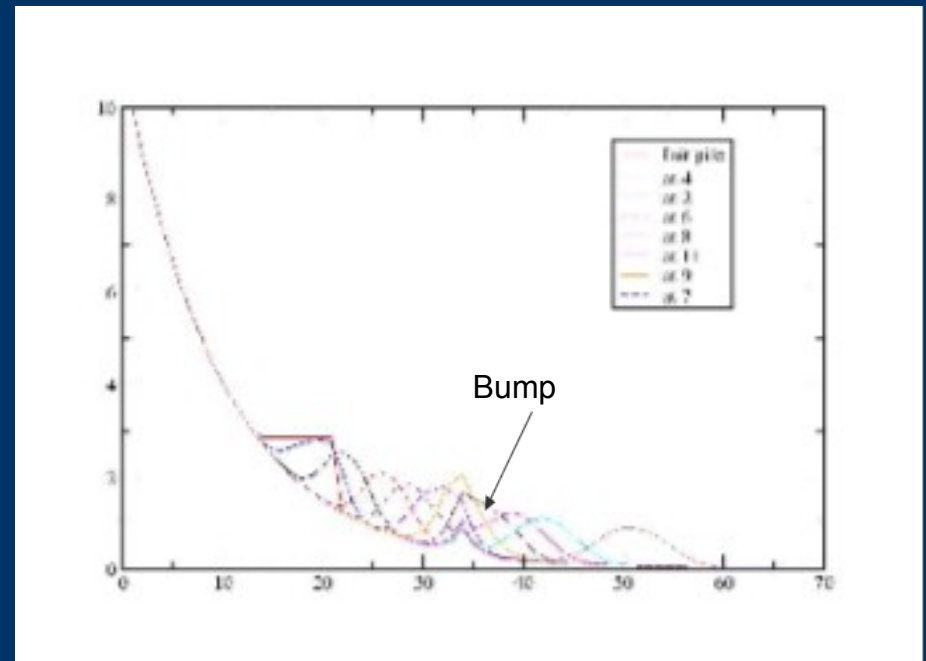


## Dam break

## Exponentially decaying topography with bump



Big bump



Small bump

# Equation system for bi-phase flows

Conservation laws, assuming a fluid saturated mixture of solids and water:

Mass:

$$\partial_t \rho^s \varphi + \nabla \cdot (\rho^s \varphi v_i) = 0$$

$$\partial_t \rho^f (1 - \varphi) + \nabla \cdot (\rho^f (1 - \varphi) u_i) = 0$$

Momentum:

$$\partial_t (\rho^s \varphi v_i) + \nabla \cdot (\rho^s \varphi v_i v_j) = \nabla \cdot T^s + f^s + \rho^s \varphi g$$

$$\partial_t (\rho^f (1 - \varphi) u_i) + \nabla \cdot (\rho^f (1 - \varphi) u_i u_j) = \nabla \cdot T^f + f^f + \rho^f (1 - \varphi) g$$

# Interaction forces between phases

$$f^s = -\varphi \nabla T^f + D(v - u)$$

$$f^f = -f^s$$

Where the drag  $D$  is approached through the empirical relationship:

$$D = \frac{3}{4} C_d \frac{\rho^f |v - u|}{f^2 d} \varphi (1 - \varphi)$$

and

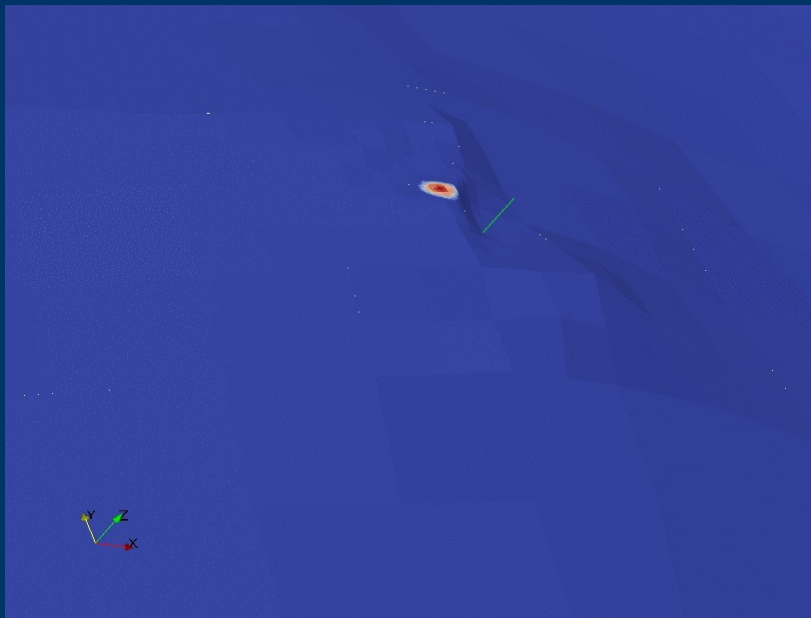
$$f = f(R_e)$$

# ***Titan2phase implementation:***

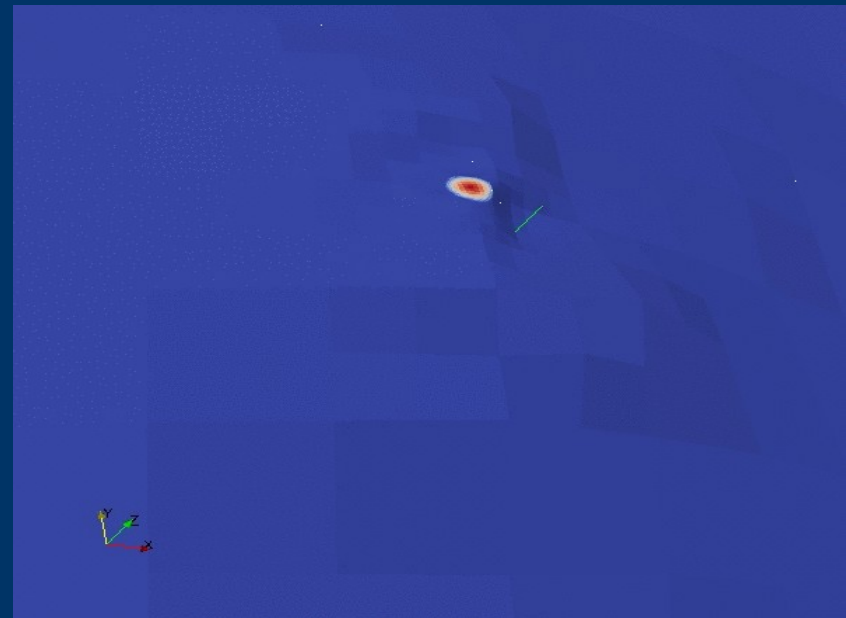
- Full two phase equation system implemented into the Titan2D frame.**
- Eigen dependent time steps of the order of  $O(10^{-4})$**
- Presence of a fluid phase reduces the bulk bed friction in a natural way.**
- Currently, constant volumetric fraction of solids assumed.**



# *Titan2D and Titan2phase*



Titan2D



Titan2phase





# *Next steps:*

- Allowing change in particle concentration in time
- Validation
- Corrections and debugging
- Validation
- Release of a beta-version
- Web distribution as free software
- Implementation on VHUB



# Equation System

## Equation system

Mass conservation

$$\delta_t \varphi_s + \frac{\delta(\varphi_s v_x)}{\delta x_1} = 0$$

$$\delta_t \varphi_f + \frac{\delta(\varphi_f u_x)}{\delta x_1} = 0$$

Momentum balance:

$$\delta_t(\varphi_s v_x) + v_j \frac{\delta(\varphi_s v_x)}{\delta x_j} = \frac{1}{\rho_s} (\nabla \cdot T_s + f_s) + \varphi_s g_x$$

$$\delta_t(\varphi_f u_x) + u_j \frac{\delta(\varphi_f u_x)}{\delta x_j} = \frac{1}{\rho_f} (\nabla \cdot T_f + f_f) + \varphi_f g_x$$

where

$$f_s = p_f \frac{\delta \varphi_s}{\delta x_1} + D \Delta u$$

$$f_f = p_f \frac{\delta \varphi_f}{\delta x_1} - D \Delta u$$

where  $\Delta u = (v - u)$  and the interface drag coefficient  $D$  is either:

$$D = \frac{\varphi_s \varphi_f^2 \Delta \rho}{\tau_s}, \text{ where the relaxation time } \tau_s = w_{hs} / g, \text{ for particles}$$

at the Stokes regime, or the empirical approach (Dobran, 2001):

$$D = \frac{3}{4} C_d \frac{\varphi_s \rho_f |\Delta u|}{d \varphi_f^2}, \text{ for other cases}$$

Constitutive equations (after depth averaging):

For solid phase, the Coulomb rule is used (Savage and Hutter (1989), Pitman and Lee (2005)):

$$T_{s_{ij}} = \alpha_{ij} T_{s_{zz}}$$

where

$$T_{s_{zz}} = \varphi_s g h$$

and (Petra et al. (2005)):

$$\alpha_{ii} = k_{ap}$$

$$\alpha_{jz} = -\frac{v_j}{\sqrt{v_x^2 + v_y^2}} \left(1 + \frac{v_j}{r \beta}\right)$$

$$\alpha_{xy} = -\text{sgn}(\partial v_x / \partial y) k_{ap} \sin \phi_{int}$$

For the fluid phase:

$$T_{f_{ij}} = \alpha_{ij} T_{f_{zz}}, \text{ where}$$

$$T_{f_{zz}} = \varphi_f g h$$

and

$$\alpha_{ii} = 1$$

$$\alpha_{jz} = -\frac{f_j}{8 u_x^2}$$

$$\alpha_{xy} = 0.5$$